

Gravitational Solitons on Einstein-Rosen Metrics

Yaron Hadad

Based on the work

‘Gravitational Solitons on Diagonal Background Metrics and on Einstein-Rosen
Background’ with V. E. Zakharov (in process)

The Theory of Relativity

The Equivalence Principle: gravity affects all bodies in the same way, independently of the body's composition.

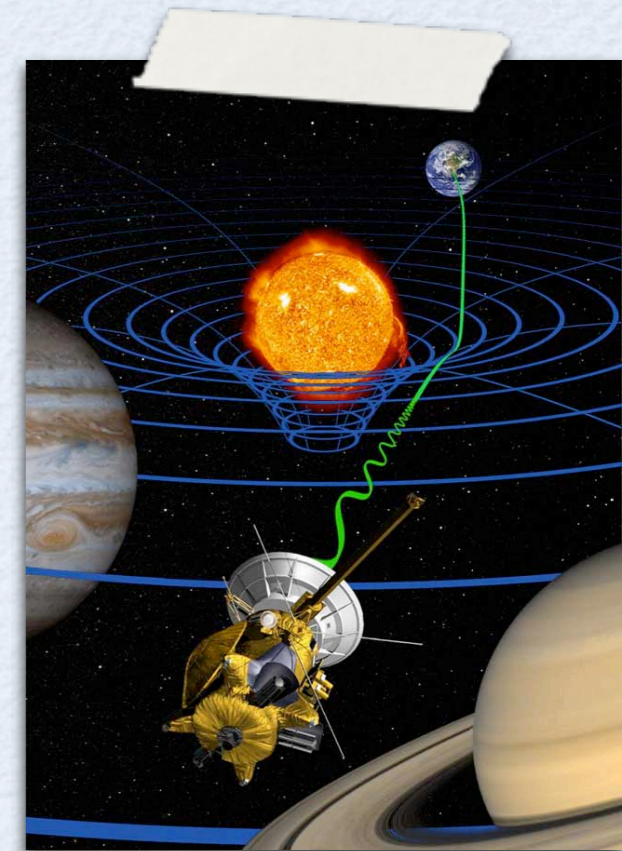
The dynamic field is the spacetime metric: $g_{\mu\nu}$

Spacetime curves in the presence of matter according to Einstein's equation:

$$\underbrace{R_{\mu\nu}}_{\text{Ricci Curvature}} - \frac{1}{2} \underbrace{g_{\mu\nu} R}_{\text{Scalar Curvature}} = (8\pi G) \underbrace{T_{\mu\nu}}_{\text{Stress-Energy Tensor}} \quad (c = 1)$$

In the absence of matter we get Einstein's vacuum equation:

$$R_{\mu\nu} = 0$$



The Setup

The spacetime metric in matrix form:

$$g_{\mu\nu} = \begin{pmatrix} -f & 0 & 0 & 0 \\ 0 & g_{11} & g_{12} & 0 \\ 0 & g_{21} & g_{22} & 0 \\ 0 & 0 & 0 & f \end{pmatrix}$$

We get the spacetime interval:

$$-ds^2 = f(-dt^2 + dz^2) + g_{ab}dx^a dx^b \quad (a, b = 1, 2)$$

Physical assumptions:

$$g_{13} = g_{23} = 0$$

$$f = f(t, z)$$

$$g_{ab} = g_{ab}(t, z)$$

Includes Schwarzschild metric, Kerr metric, Einstein-Rosen metric, Bianchi metrics, Kasner metric and many more...

Einstein's Equation

Use 'light-cone' coordinates: $t = \zeta - \eta$ $z = \zeta + \eta$

Einstein's vacuum equation:

$$R_{\mu\nu} = 0$$

$$R_{ab} = 0$$

$$\begin{aligned} R_{00} + R_{33} &= 0 \\ R_{03} &= 0 \end{aligned}$$

$$\underbrace{(\alpha g_{,\zeta} g^{-1})_{,\eta}}_{-A} + \underbrace{(\alpha g_{,\eta} g^{-1})_{,\zeta}}_B = 0$$

$$(\ln f)_{,\zeta} = \frac{(\ln \alpha)_{,\zeta\zeta}}{(\ln \alpha)_{,\zeta}} + \frac{1}{4\alpha\alpha_{,\zeta}} \text{Tr} (A^2)$$

$$(\ln f)_{,\eta} = \frac{(\ln \alpha)_{,\eta\eta}}{(\ln \alpha)_{,\eta}} + \frac{1}{4\alpha\alpha_{,\eta}} \text{Tr} (B^2)$$

$$\det (g) = \alpha^2$$

$$\alpha_{,\zeta\eta} = 0 \implies \alpha(\zeta, \eta) = a(\zeta) + b(\eta) \quad \beta(\zeta, \eta) = a(\zeta) - b(\eta)$$

Belinski & Zakharov (1978) found a Lax pair for the $g(\zeta, \eta)$ eq.

The Inverse Transform

1

Take a particular solution $g_0(\zeta, \eta)$

Can skip this step when g_0 is diagonal

2

$$\begin{aligned} D_1 \psi_0 &= \frac{A_0}{\lambda - \alpha} \psi_0 \\ D_2 \psi_0 &= \frac{B_0}{\lambda + \alpha} \psi_0 \end{aligned} \implies \text{Get } \psi_0$$

3

Define the dressing matrix:

$$\chi = I + \sum_{k=1}^n \left(\frac{R_k}{\lambda - \mu_k} + \frac{\bar{R}_k}{\lambda - \bar{\mu}_k} \right)$$

$$\mu_k = \omega_k - \beta \pm \sqrt{(\omega_k - \beta)^2 - \alpha^2}$$

R_k can be expressed in terms of ψ_0

4

Get a solitonic solution!

$$g = \chi(0) g_0$$

Diagonal Metrics

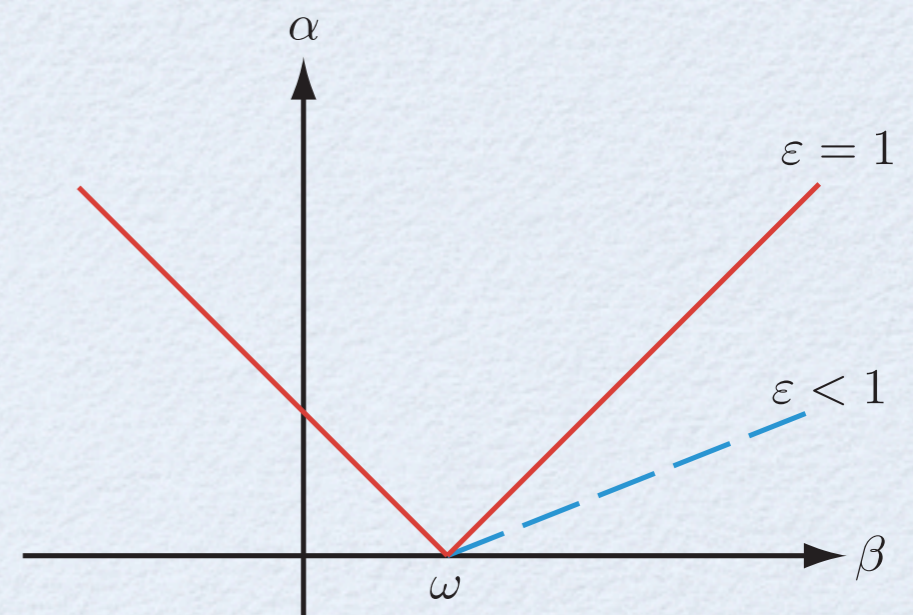
$$g_0 = \text{diag}(\alpha e^{2s}, \alpha e^{-2s})$$

\implies n-soliton solution can be written explicitly in this case, but will depend on the functions ρ_k :

Cannot always be integrated analytically.

$$\left\{ \begin{array}{l} \rho_{k,\zeta} = \frac{\alpha + \mu_k}{\alpha - \mu_k} 2s_{,\zeta} \\ \rho_{k,\eta} = \frac{\alpha - \mu_k}{\alpha + \mu_k} 2s_{,\eta} \end{array} \right.$$

$$\begin{aligned} \mu_k &= \omega_k - \beta \pm \sqrt{(\omega_k - \beta)^2 - \alpha^2} \\ &= (\beta - \omega_k) (-1 \pm \sqrt{1 - \varepsilon_k^2}) \end{aligned}$$



$$\varepsilon_k = \frac{\alpha}{\beta - \omega_k}$$

Write: $\rho_k = \rho_k^{(0)} + \rho_k^{(1)} \varepsilon_k + \rho_k^{(2)} \varepsilon_k^2 + \dots$

\implies Get a recursive relation for ρ_k in terms of s only!

One-Soliton on Einstein-Rosen Background

Einstein-Rosen metric:

$$\begin{aligned}
 t &\mapsto r & g_0 &= \text{diag}(\alpha e^{2s}, \alpha e^{-2s}) \\
 z &\mapsto t & \alpha(r, t) &= r \\
 x &\mapsto i\theta & s(r, t) &= J_0(r) \sin(t) \\
 y &\mapsto iz & &
 \end{aligned}$$

$$\varepsilon_k = \frac{r}{t - \omega_k}$$

$$\begin{aligned}
 \rho_k^{(0)}(r, t) &= 2J_0(r) \sin(t) \\
 \rho_k^{(1)}(r, t) &= -2J_1(r) \cos(t) \\
 &\vdots
 \end{aligned}$$

One-soliton solution:

$$g = \frac{1}{2\mu_1 \cosh \rho_1} \begin{pmatrix} (\mu_1^2 e^{\rho_1} + \alpha^2 e^{-\rho_1}) e^{2s} & \alpha^2 - \mu_1^2 \\ \alpha^2 - \mu_1^2 & (\alpha^2 e^{\rho_1} + \mu_1^2 e^{-\rho_1}) e^{-2s} \end{pmatrix}$$

$$f = f_0 \sqrt{\alpha \mu_1} \cosh \left(\frac{\rho_1}{\alpha^2 - \mu_1^2} \right)$$

In this case the functions ρ_k are given by:

$$\rho_k(r, t) = 2J_0(r) \sin(t) - 2J_1(r) \cos(t) \varepsilon_k + \left[J_0(r) \sin(t) - \frac{2J_1(r) \sin(t)}{r} \right] \varepsilon_k^2 + O(\varepsilon_k^3)$$

Pretty Soliton

$$g = \frac{1}{2\mu_1 \cosh \rho_1} \begin{pmatrix} (\mu_1^2 e^{\rho_1} + \alpha^2 e^{-\rho_1}) e^{2s} & \alpha^2 - \mu_1^2 \\ \alpha^2 - \mu_1^2 & (\alpha^2 e^{\rho_1} + \mu_1^2 e^{-\rho_1}) e^{-2s} \end{pmatrix}$$

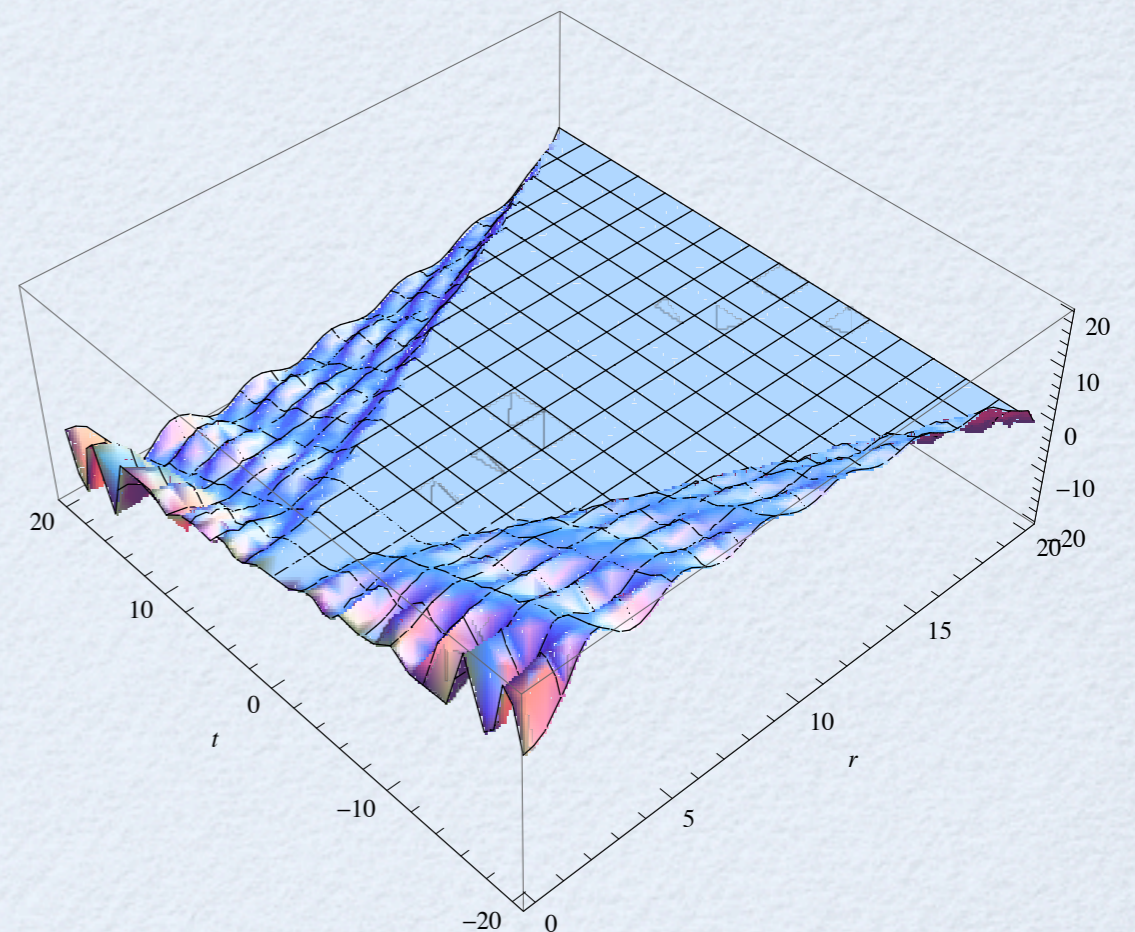
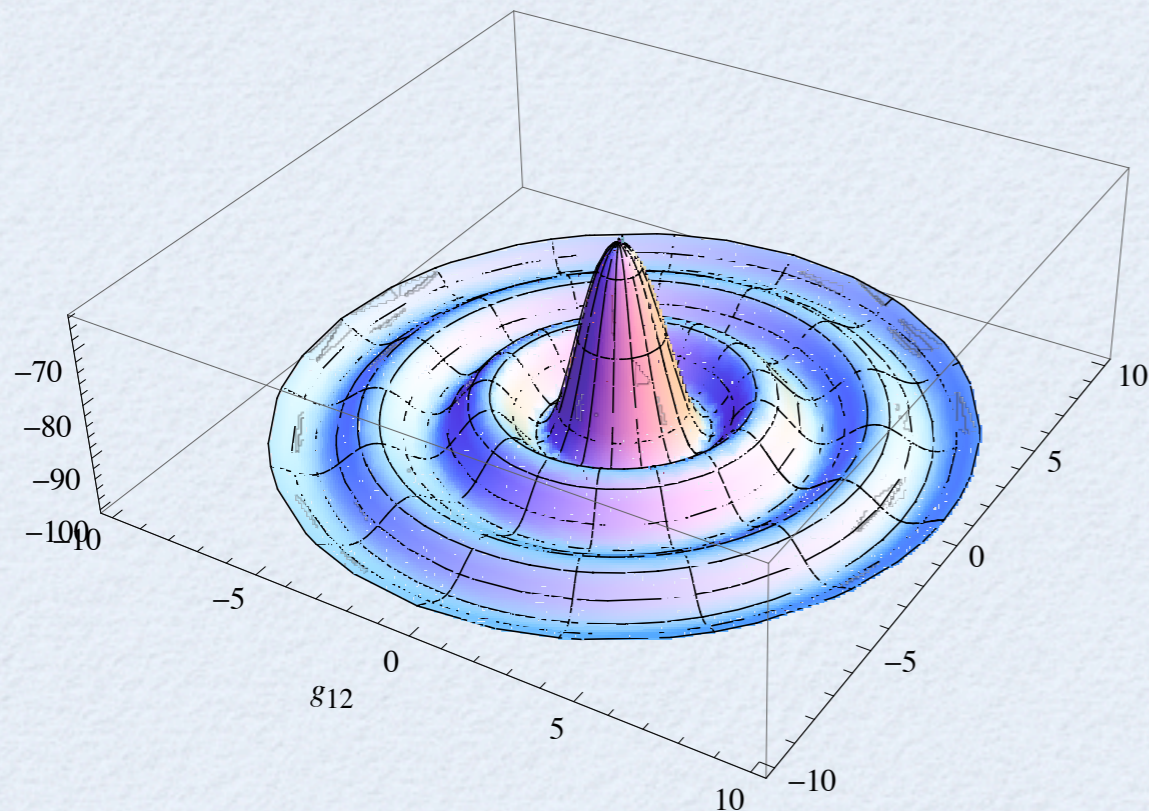
$$\omega_1 = 0$$

What happens in the limit $\varepsilon \rightarrow 1$?

$$g \rightarrow g_0$$

$$0 \leq r \leq 10 \quad t = 100$$

\implies Solution can be extended.



Two Solitons

In this case there are two integration constants: $\omega_1, \omega_2 \in \mathbb{R}$

The light-cones divided spacetime into 6 regions:

I, VI : The two-soliton region

III : The two-soliton 'interaction' region

II, V : One-Soliton regions

IV : Background metric

$$\varepsilon_k = \frac{r}{t - \omega_k}$$

